

# Low-Frequency Variability in a Low-Order Shallow Water Atmospheric Model

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## Abstract

Low-frequency variability is investigated in a low-order atmospheric model derived from the 2-layer shallow water equations on a  $\beta$ -plane channel with bottom topography.

### • Low-frequency variability in our model:

- variability on time scales from 10 to 200 days on a zonal wavenumber 3 (Benzi & Speranza 1989);
- irregular weakening and amplification of non-propagating planetary waves;
- related to strange attractors through bifurcations of periodic orbits;
- dominant time scales and spatial patterns: inherited from periodic orbits;
- persistent in parameter space.
- New scenario: available theories (Charney & DeVore 1979, Legras & Ghil 1985, Crommelin et al. 2004) rely on multiple equilibria and unrealistic zonal wind-speeds.

## Model

- 46-dimensional; obtained from SW equations by Galerkin projection on wavenumbers 0,3 (zonal) and 0,1,2 (meridional).
- Forced by zonal wind profile with magnitude  $U_0$ .
- Planetary-scale channel ( $29000 \times 2500$  km).
- Pressure in each layer

$$p_1 = \rho_1 g(h_1 + h_2 + h_b - z)$$

$$p_2 = \rho_2 g(h_1 + \rho_2 g(h_2 + h_b - z))$$

### • Dynamical equation in each layer

$$u_t + uu_x + vv_y = -\frac{p_x}{\rho} + (f + \beta y)v$$

$$v_t + uv_x + vv_y = -\frac{p_y}{\rho} - (f + \beta y)u$$

$$h_t + uh_x + vh_y = -h(u_x + v_y)$$

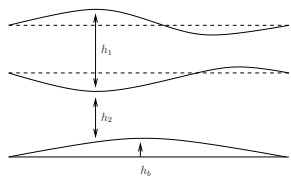


Figure 1: Setup of the 2-layer shallow water model. Each layer has a constant density and variable thickness.

## Bifurcations

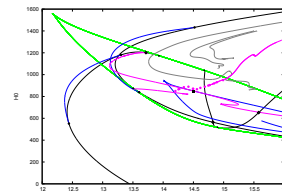


Figure 2: Bifurcation curves in the  $(U_0, h_0)$ -plane: saddle-node bifurcation of equilibria (green), Hopf bifurcation of equilibria (black), saddle-node bifurcation of periodic orbits (blue), period doubling bifurcations (grey), and Neimark-Sacker bifurcations (magenta).

### • Bifurcations of equilibria and periodic orbits:

- control parameters: magnitude of zonal wind forcing ( $U_0$ ) and topographic height ( $h_0$ );
- there exist two different branches of periodic orbits:  $h_0 \leq 950$ : period 10 days, traveling wave
- $h_0 \geq 950$ : period 50 days, non-propagating wave
- codimension-2 points: Hopf-saddle-node of equilibria (triangles) and periodic orbits (square), generalized Hopf (circle), double Hopf (diamond);

### • Routes to chaos:

- period-doubling cascade;
- breakdown of a 2-torus attractor after a Neimark-Sacker bifurcation;
- saddle-node bifurcations of periodic orbits.

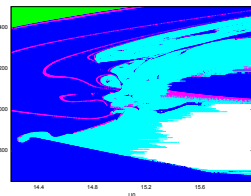


Figure 3: Lyapunov diagram of attractors. Equilibria (green), periodic attractors (blue), 2-torus attractors (magenta), and strange attractors (cyan).

### • Occurrence of attractors:

- strange attractors appear when  $14 \leq U_0 \leq 16$  and  $650 \leq h_0 \leq 1500$ ;
- islands of periodic behaviour in the middle of chaos.

## Strange attractor I

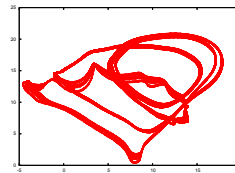


Figure 4: Strange attractor after period doubling cascade ( $U_0 = 15, h_0 = 1200$ ).

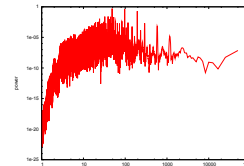


Figure 5: Power spectrum: dominant time scale approximately 50 days; peaks at multiples of 50 days due to period doublings.

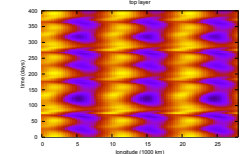


Figure 6: Hovmöller diagram of vorticity along a line of constant latitude: irregular weakening and amplification of vorticity field; non-propagating character; variability stronger in upper layer.

## Strange attractor II

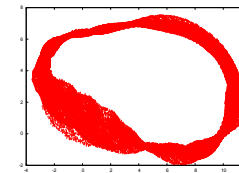


Figure 7: Strange attractor after breakdown of a 2-torus attractor ( $U_0 = 14.78, h_0 = 900$ ).

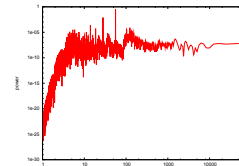


Figure 8: Power spectrum: dominant time scales are 20 and 60 days; two peaks are inherited from the formerly existing 2-torus.

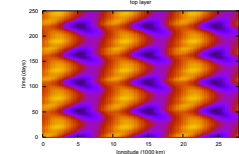


Figure 9: Hovmöller diagram of vorticity along a line of constant latitude: irregular weakening and amplification of vorticity field; non-propagating character; time scales different for both layers.

## Conclusions

### • Summary of dynamical scenario:

- Hopf bifurcations give rise to two different branches of periodic orbits;
- periodic attractors bifurcate into strange attractors via three different routes to chaos;
- strange attractors are related to irregular weakening and amplification of vorticity field;

### • The dynamical mechanisms in our model are very different from the available theories (Legras & Ghil 1985, Crommelin et al. 2004).

### • Future work:

- focus on the dynamics around the codimension-2 points in the parameter plane;
- determination of the boundaries of the periodic islands within the chaotic zone;
- investigate whether the current dynamical scenarios persist when the number of retained modes in the Galerkin projection is increased.

## References

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